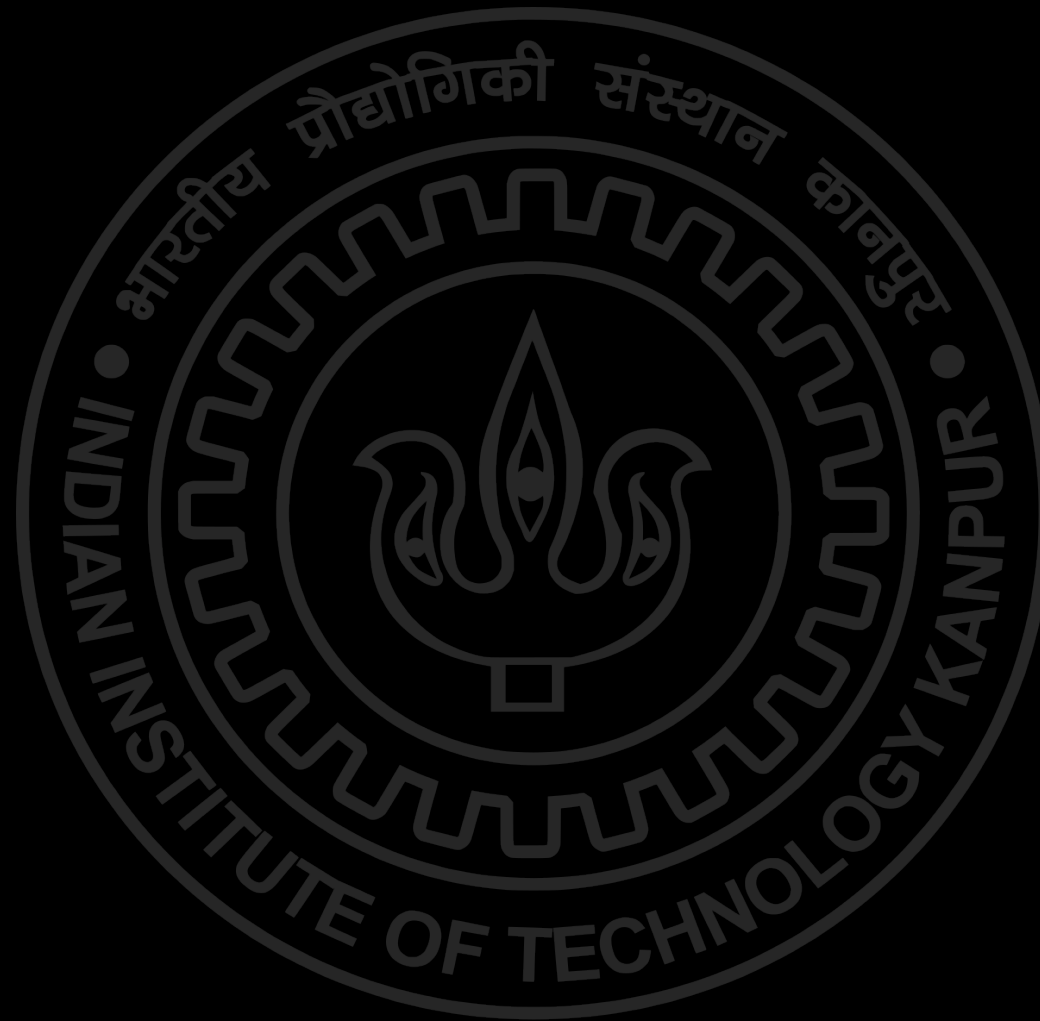


Spectral Method for Solving PDEs

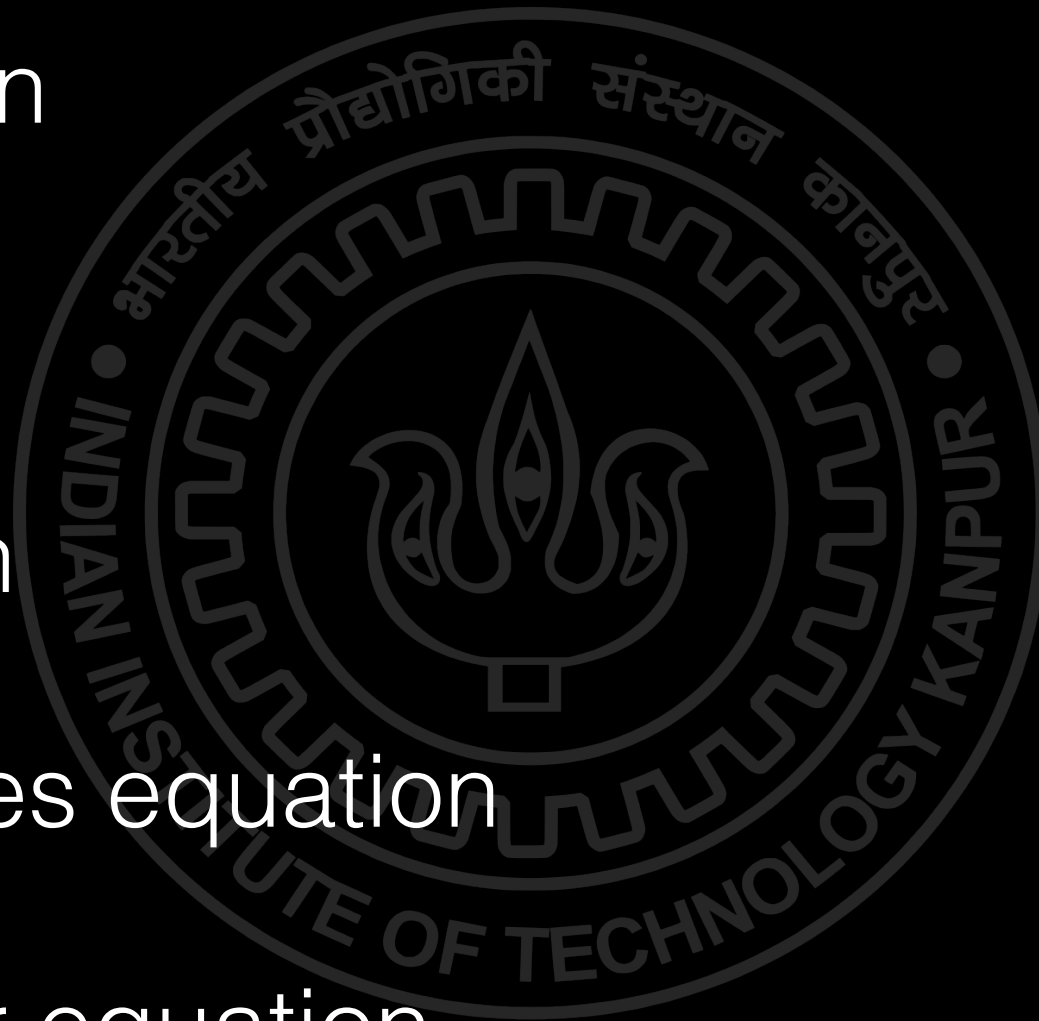
Mahendra Verma

$$\frac{\partial \phi}{\partial t} = f\left(\phi, \frac{\partial \phi}{\partial x}, \frac{\partial^2 \phi}{\partial x^2}, t\right)$$

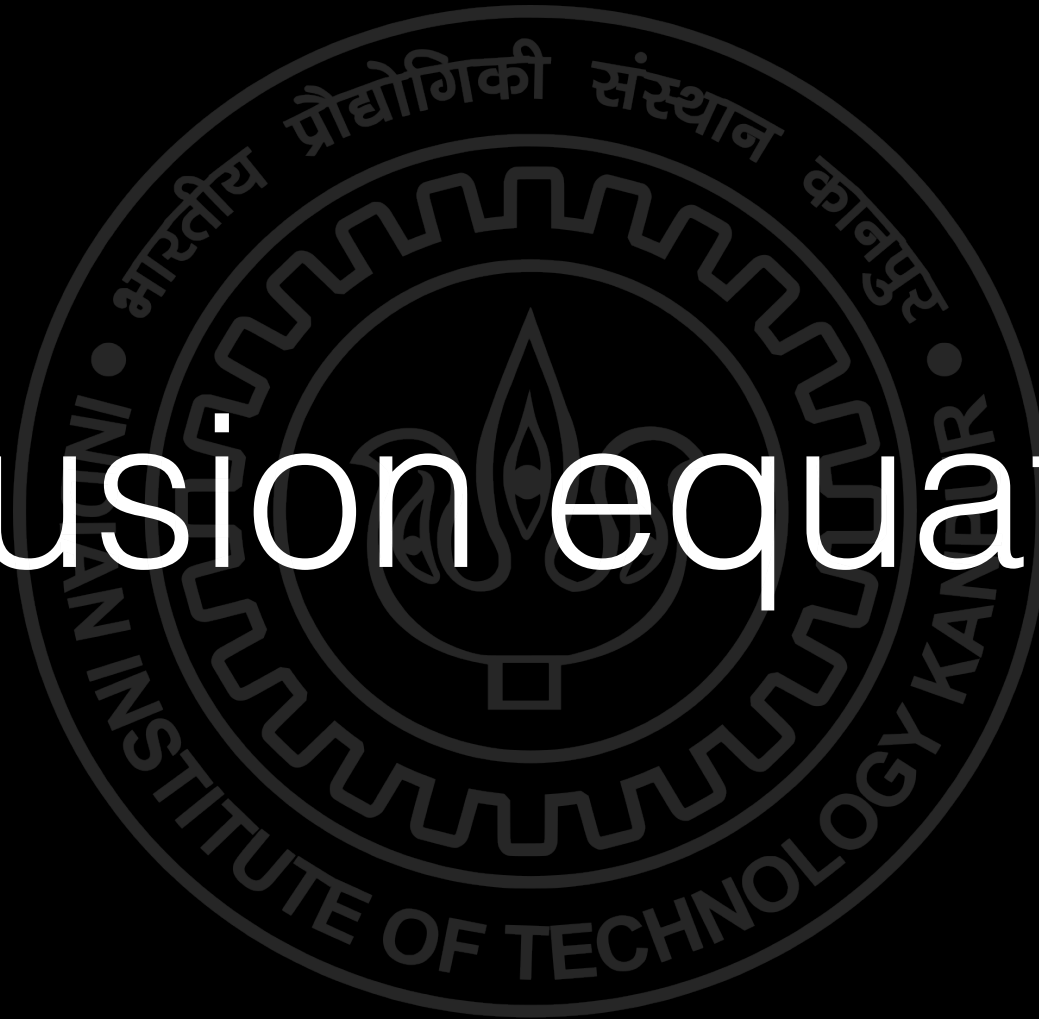


PDEs

- Diffusion eqn
- Wave eqn
- Burgers eqn
- Navier-Stokes equation
- Schrodinger equation
-



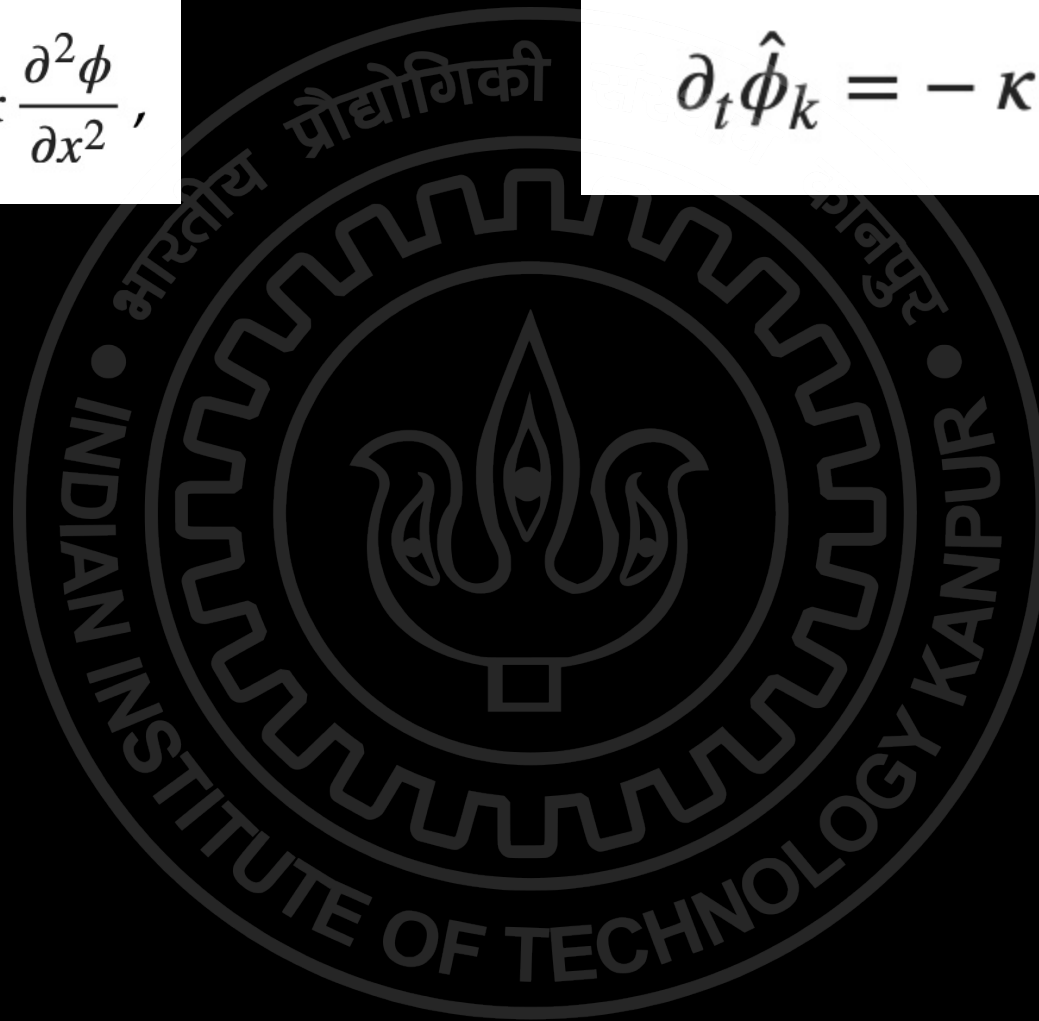
Diffusion equation



Diffusion equation

$$\partial_t \phi = \kappa \frac{\partial^2 \phi}{\partial x^2},$$

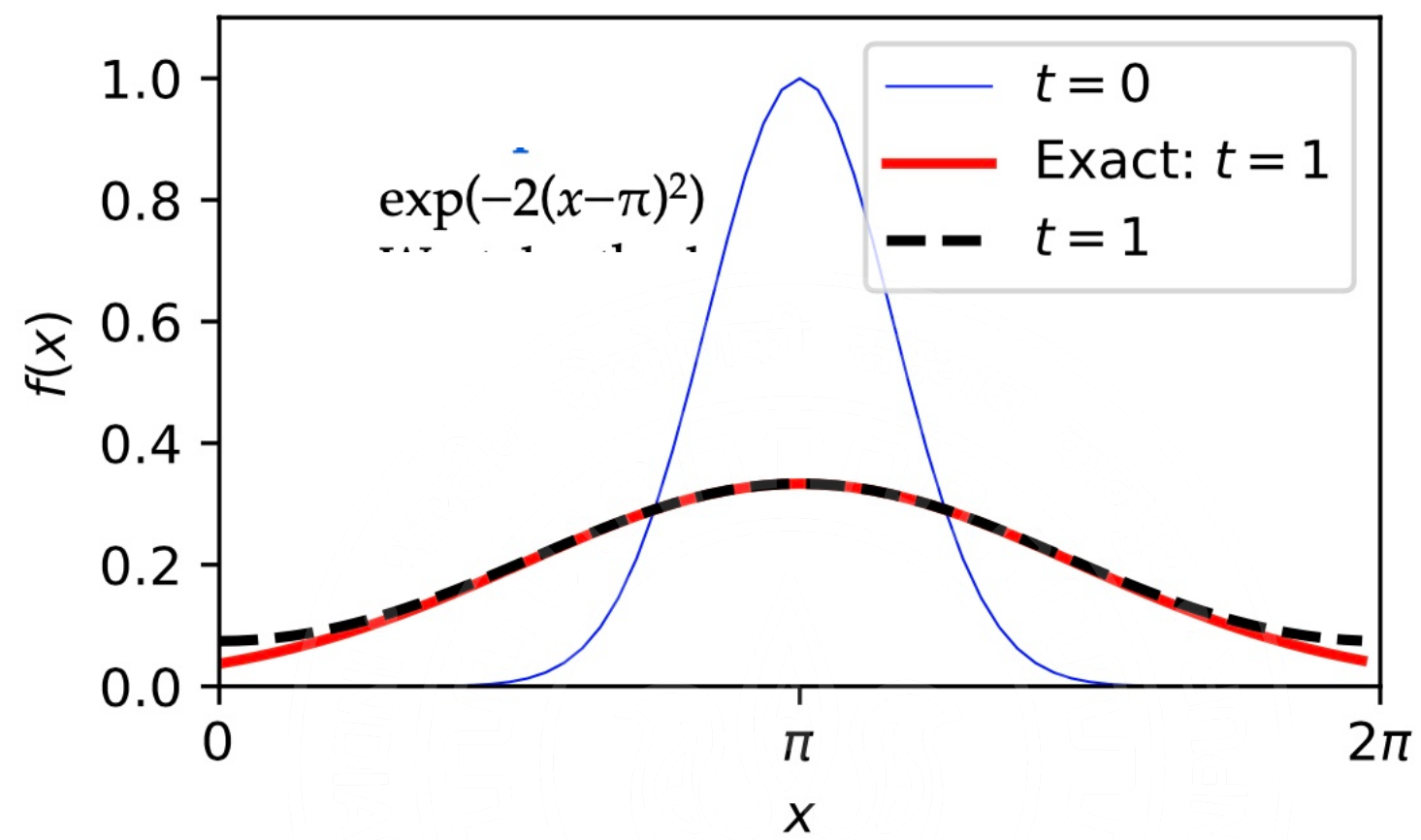
$$\partial_t \hat{\phi}_k = -\kappa k^2 \hat{\phi}_k,$$



```
kappa = 1; tf = 1
L = 2*np.pi; N = 64; h = L/N
j = np.arange(0,N); x = j*h

f = np.exp(-2*(x-pi)**2) # init cond
fk = np.fft.rfft(f,N)/N
kx = np.linspace(0, N//2, N//2+1)

# Final states
fk_t = fk*np.exp(-kappa*kx**2*tf)
f_t = np.fft.irfft(fk_t,N)*N # in Real space
```



CFL condition

$$\partial_t \hat{\phi}_k = -\kappa k^2 \hat{\phi}_k,$$

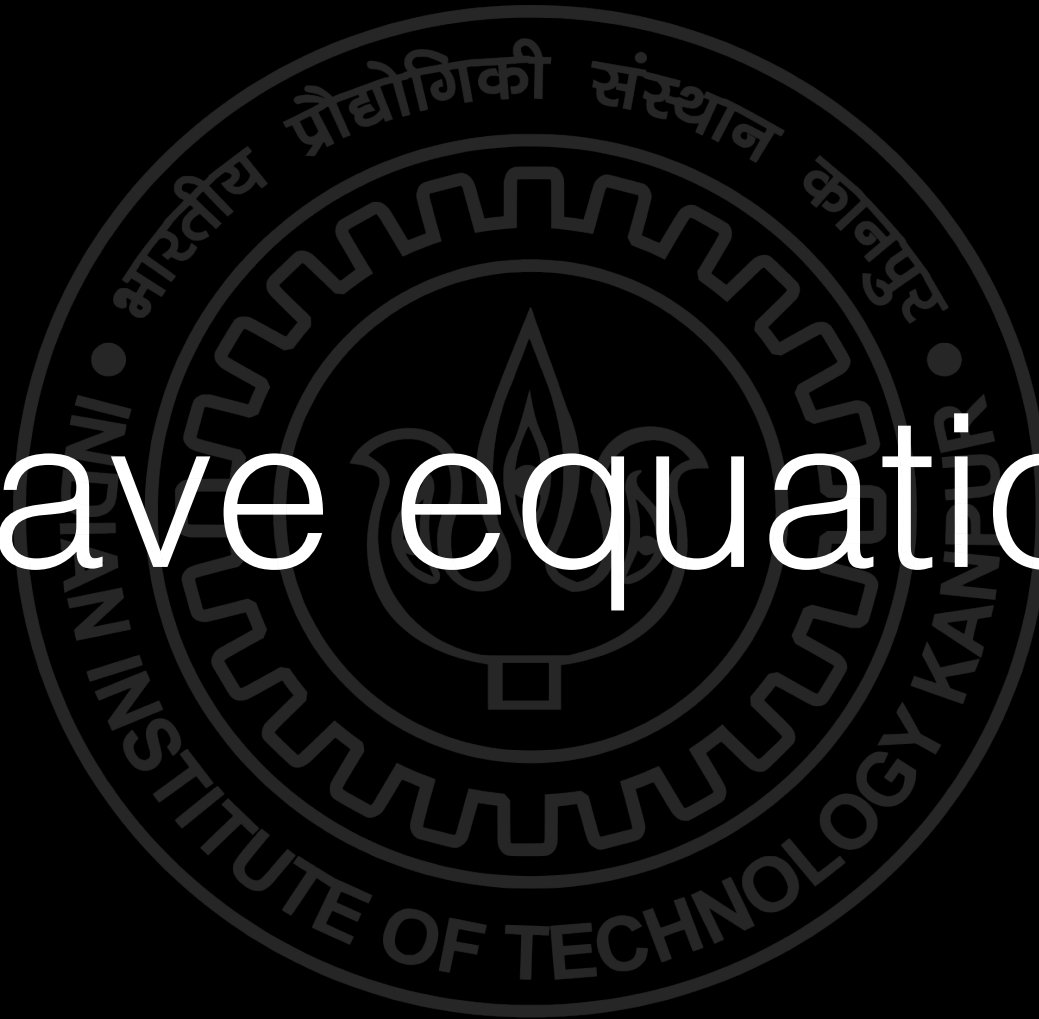
$$\hat{\phi}_k^{(n+1)} = \hat{\phi}_k^{(n)} [1 - \kappa k^2 (\Delta t)].$$

$$|1 - \kappa k^2 (\Delta t)| < 1, \quad \text{or} \quad \Delta t < \frac{2}{\kappa k^2}.$$

$$\Delta t < \frac{2}{\kappa k_{\max}^2}.$$

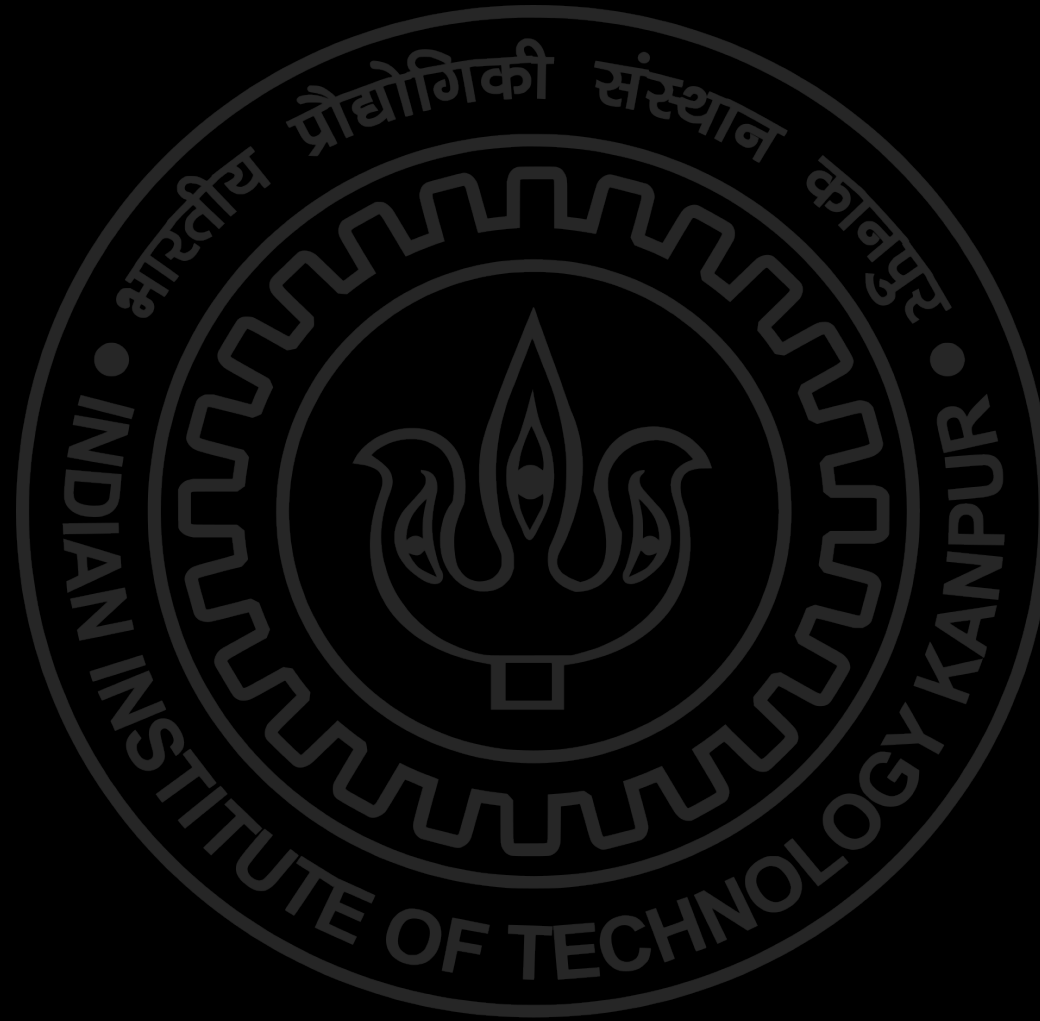
$$\Delta t < \frac{2h^2}{\kappa \pi^2}.$$

Wave equation



$$\partial_t \phi + c \partial_x \phi = 0 ,$$

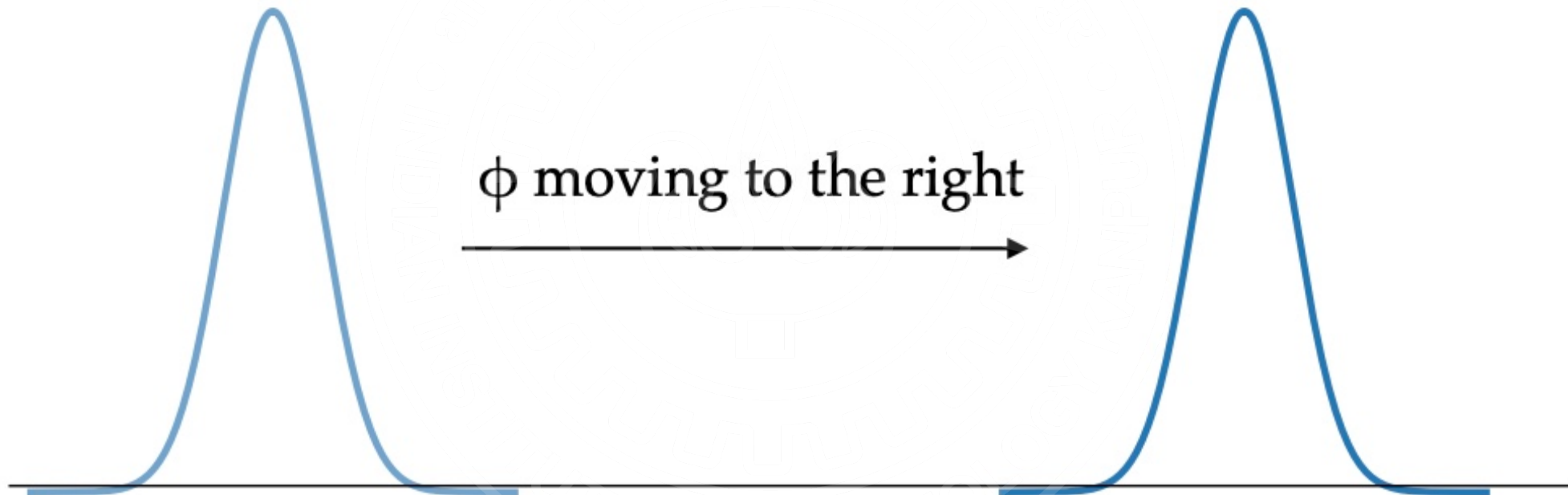
$$\frac{d}{dt} \hat{\phi}_k = -i c k \hat{\phi}_k .$$



$\phi(x, t = 0)$

$\phi(x, t)$

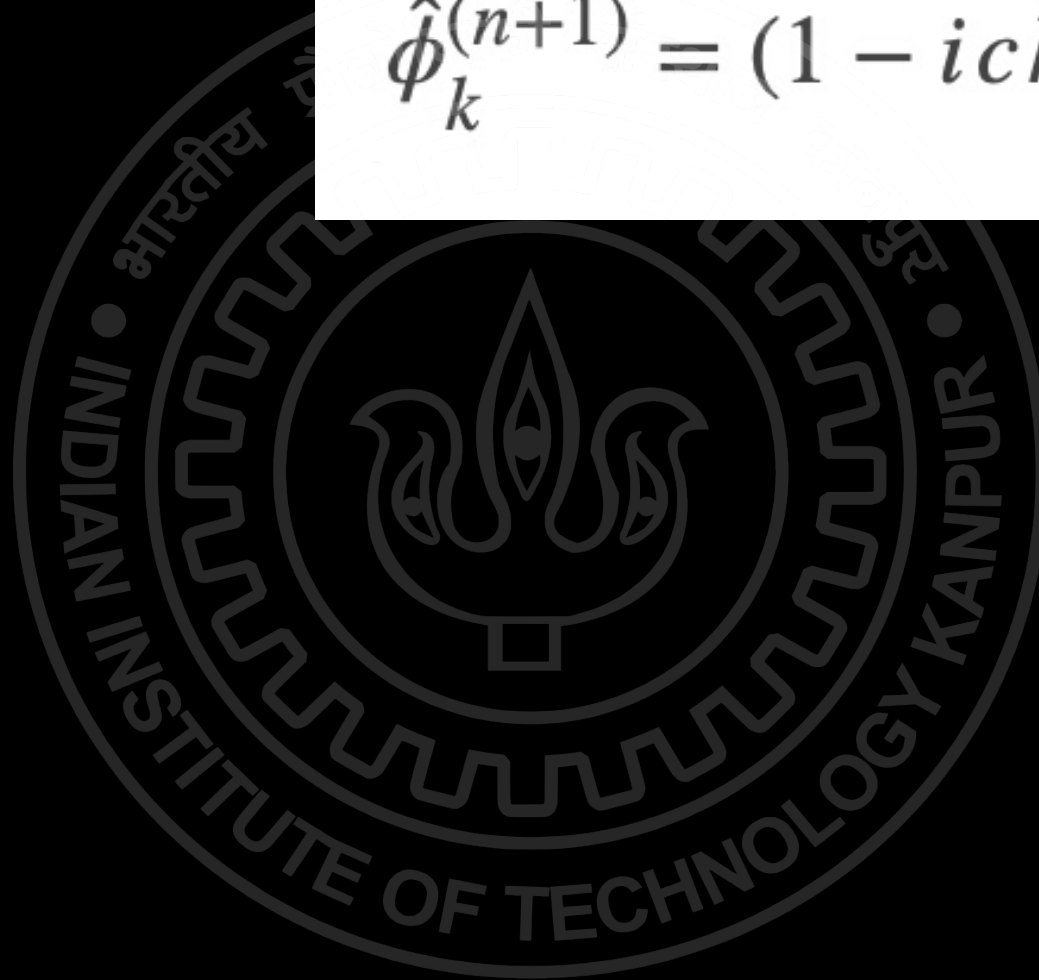
ϕ moving to the right



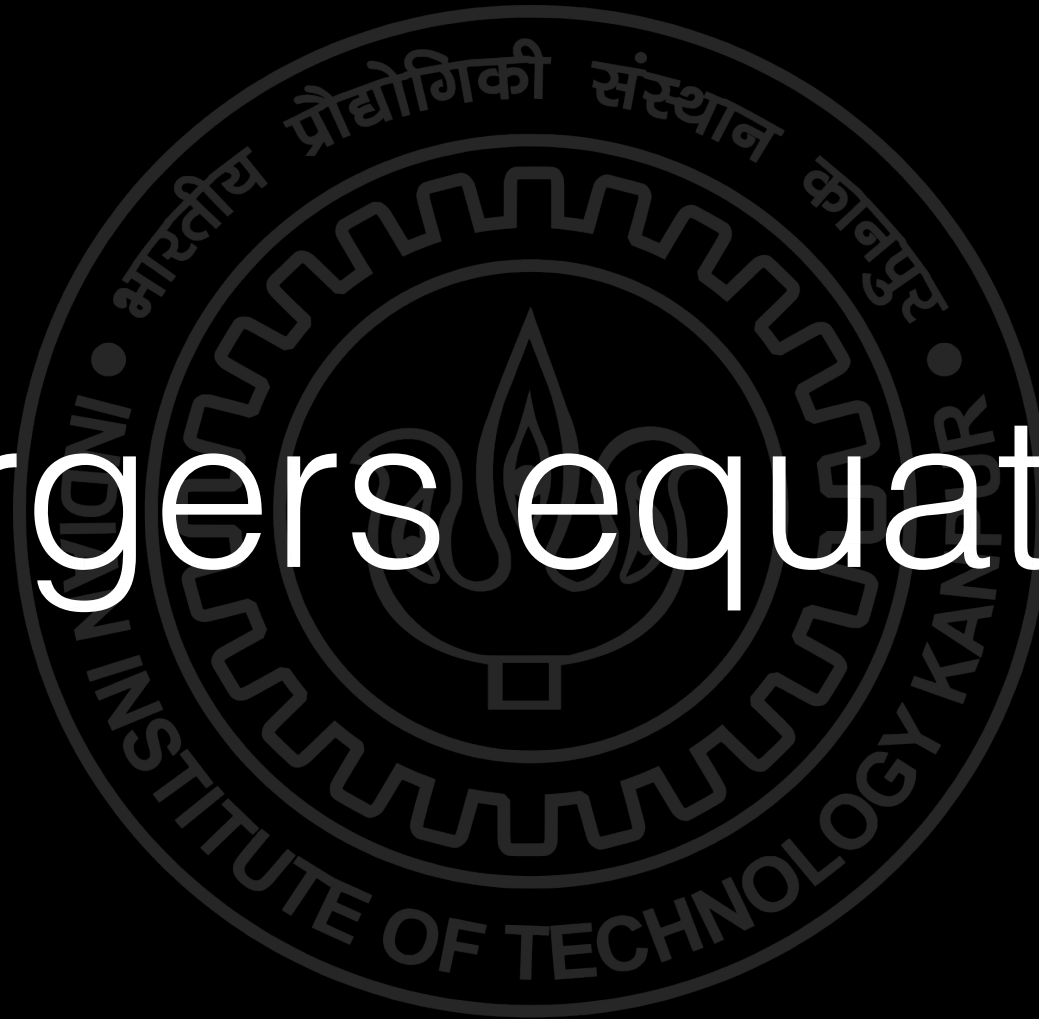
Stability

$$\frac{d}{dt}\hat{\phi}_k = -ick\hat{\phi}_k.$$

$$\hat{\phi}_k^{(n+1)} = (1 - ick(\Delta t))\hat{\phi}_k^{(n)}$$



Burgers equation



$$\partial_t u + u \partial_x u = \nu \partial_x^2 u, \text{ or } \partial_t u + \partial_x(u^2/2) = \nu \partial_x^2 u ,$$

$$\frac{d}{dt} \hat{u}_k = -i \frac{k}{2} \widehat{u^2} - \nu k^2 \hat{u}_k = -\hat{N}_k - \nu k^2 \hat{u}_k,$$

$$\hat{u}_k^{(n+1)} = \hat{u}_k^{(n)} [1 - \nu k^2 (\Delta t)] - (\Delta t) \hat{N}_k.$$

$$|1 - \nu k^2 (\Delta t)| < 1, \text{ or } \Delta t < 2/(\nu k^2).$$

Crank-Nicolson scheme

$$\frac{d}{dt}\hat{u}_k = -i\frac{k}{2}\widehat{u^2} - \nu k^2\hat{u}_k = -\hat{N}_k - \nu k^2\hat{u}_k,$$

$$\hat{u}_k^{(n+1)} = \hat{u}_k^{(n)} - \nu k^2(\Delta t)\frac{1}{2}\left(\hat{u}_k^{(n)} + \hat{u}_k^{(n+1)}\right) - (\Delta t)\hat{N}_k,$$

$$\text{or, } \hat{u}_k^{(n+1)} = \hat{u}_k^{(n)} \left[\frac{1 - \nu k^2(\Delta t)/2}{1 + \nu k^2(\Delta t)/2} \right] - (\Delta t)\hat{N}_k$$

Exponential Trick

$$\frac{d}{dt}\hat{u}_k = -i\frac{k}{2}\widehat{u^2} - \nu k^2\hat{u}_k = -\hat{N}_k - \nu k^2\hat{u}_k,$$

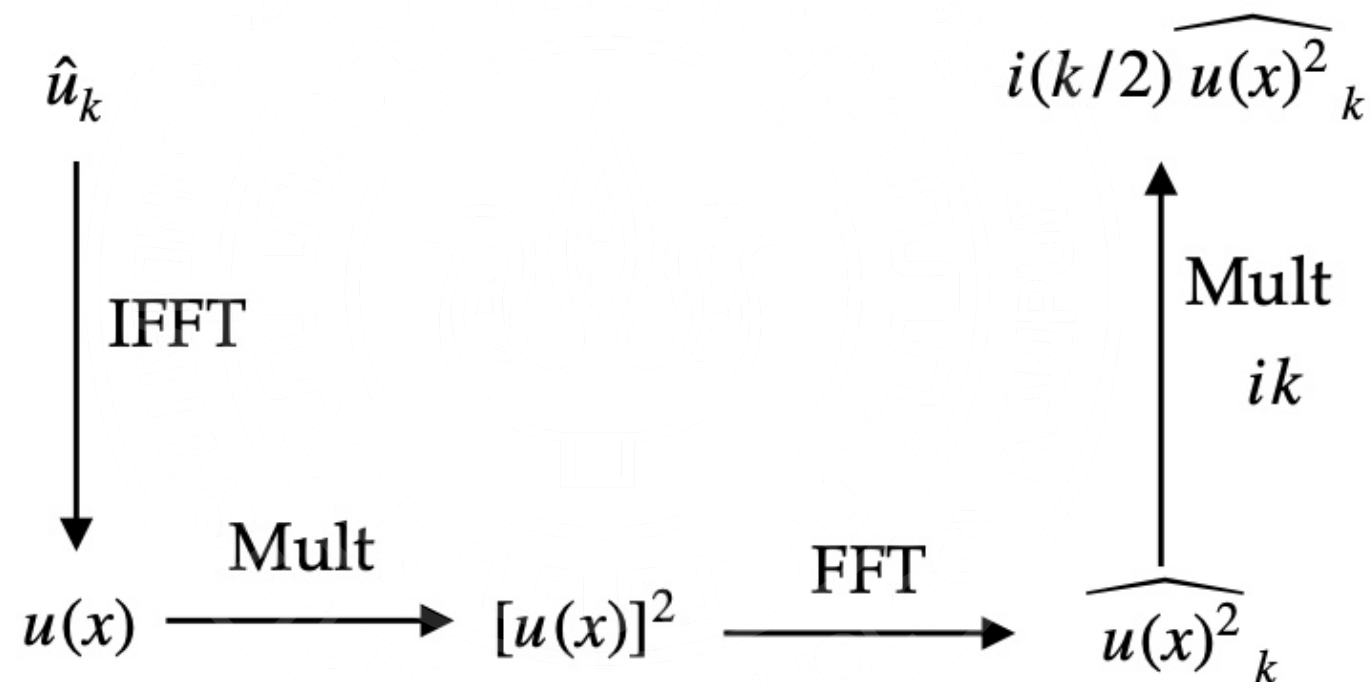
$$\hat{u}'_k = \hat{u}_k \exp(\nu k^2 t),$$

$$\frac{d}{dt}\hat{u}'_k = -\hat{N}_k \exp(\nu k^2 t).$$

An application Euler's forward scheme to the above equation yields

$$\hat{u}_k^{(n+1)} = \left[\hat{u}_k^{(n)} - (\Delta t)\hat{N}_k^{(n)} \right] \exp(-\nu k^2(\Delta t)).$$

Convolution computation



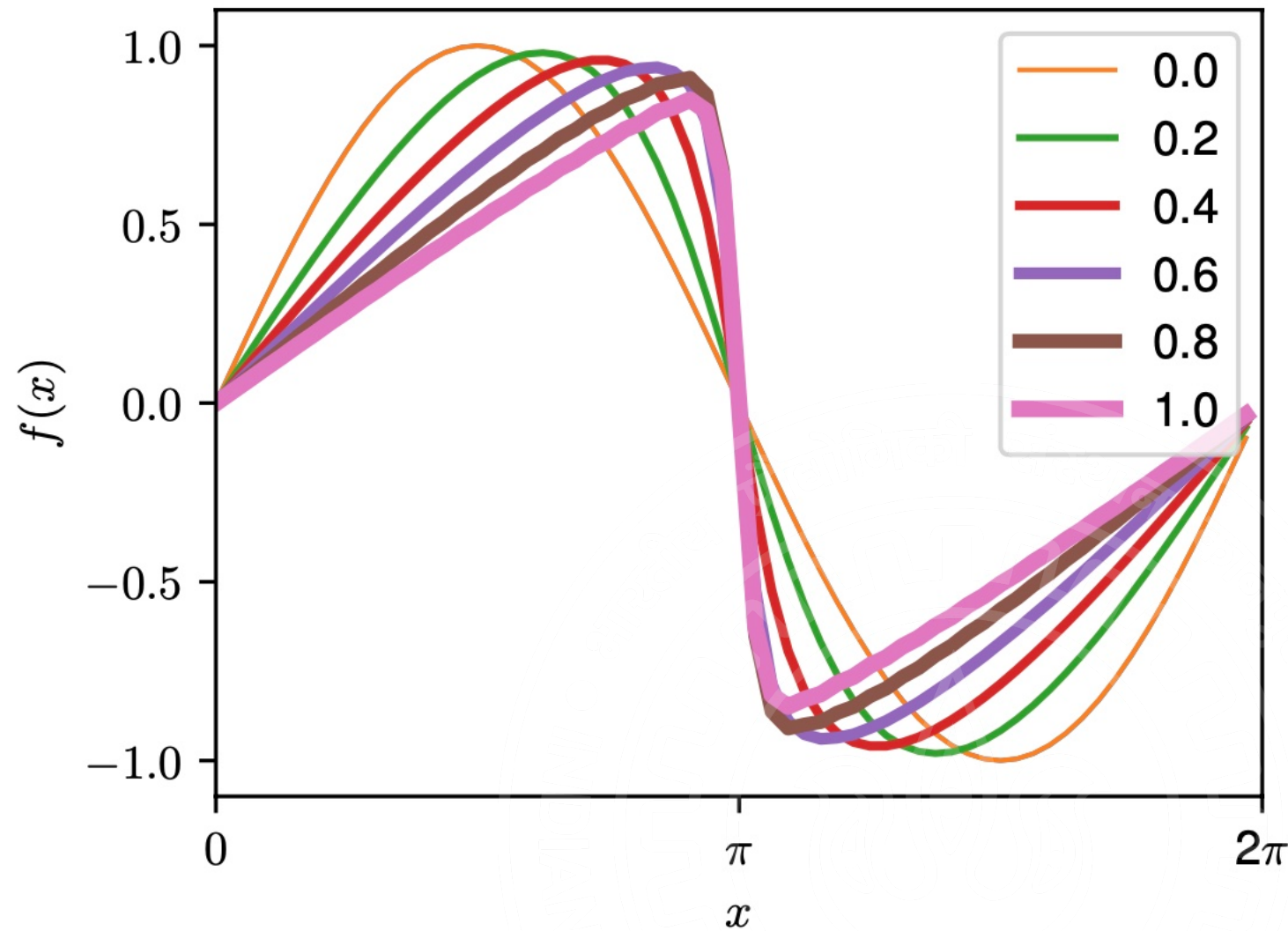
```
nu = 0.1; tf = 1; dt = 0.001; nsteps = int(tf/dt)
L = 2*np.pi; N = 64; h = L/N
j = np.arange(0,N); x = j*h

kx = np.linspace(0, N//2, N//2+1)
exp_factor_dtby2 = np.exp(-nu*kx**2*dt/2)
exp_factor = np.exp(-nu*kx**2*dt)

def comput_Nk(fk):
    f = np.fft.irfft(fk,N)*N
    f = f*f
    fk_prod = np.fft.rfft(f,N)/N
    return (1j*kx*fk_prod)

f = np.sin(x) # initiation condition
fk = np.fft.rfft(f,N)/N # FT(f)
```

```
for i in range(nsteps+2):
    Nk = comput_Nk(fk)
    fk_mid = (fk -(dt/2)*Nk)*exp_factor_dtby2
    Nk_mid = comput_Nk(fk_mid)
    fk = (fk -dt*Nk_mid)*exp_factor
```



$$(\Delta t)_{\min} = \frac{2h^2}{\nu \pi^2} = \frac{2 \times 4\pi^2}{0.1 \times 64^2 \pi^2} = \frac{10}{512} \approx 2 \times 10^{-2}.$$

On the other hand, based on the nonlinear term,

$$\Delta t < h/U_{\text{rms}} \approx 1/64 \approx 0.015.$$

$$\Delta t = 0.001,$$

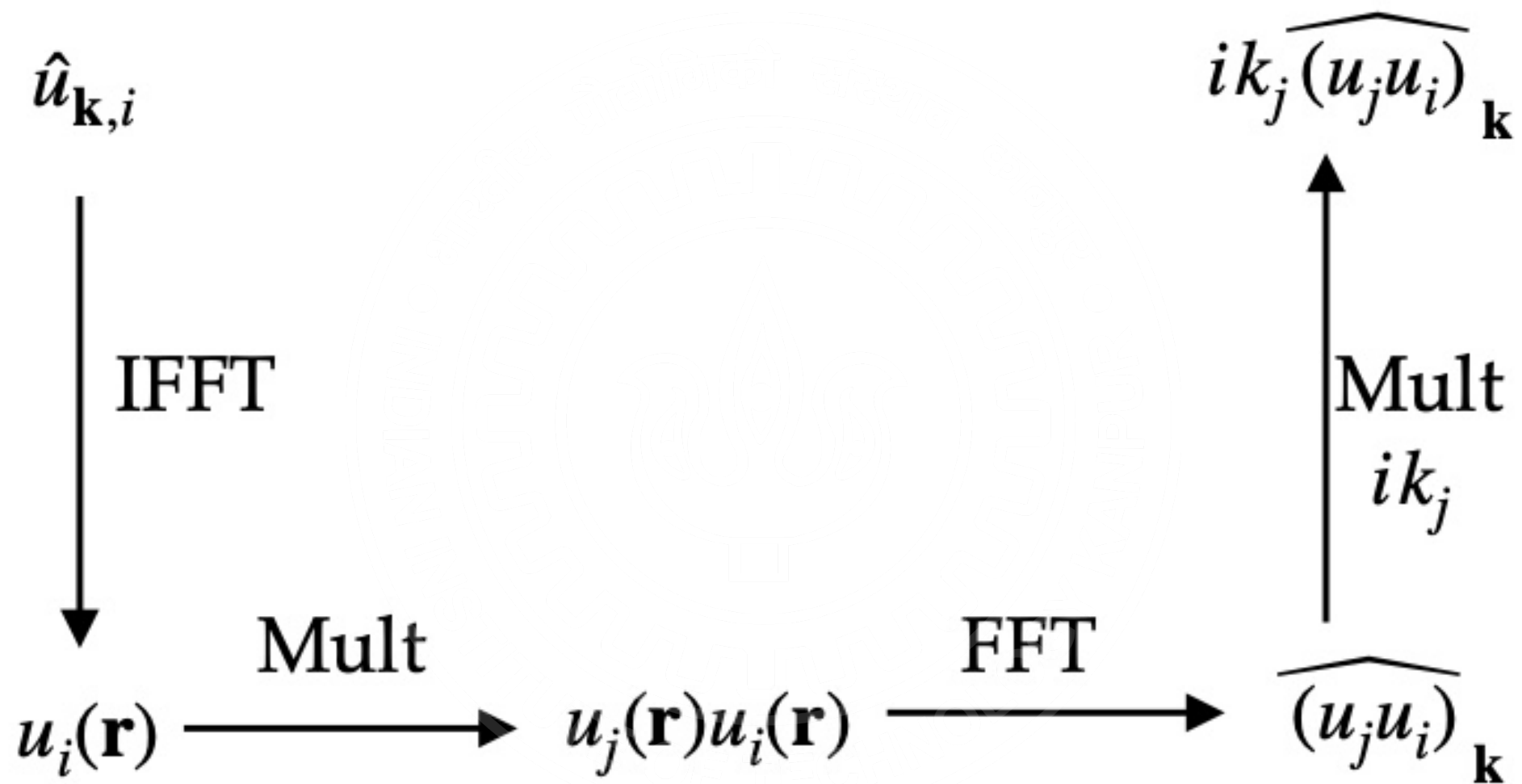
Navier-Stokes eqn

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u}, \quad \dots (80)$$

$$\nabla \cdot \mathbf{u} = 0, \quad \dots (81)$$

$$\begin{aligned} \frac{d}{dt} \hat{u}_{\mathbf{k},i} &= -i k_j \widehat{u_j(\mathbf{r}) u_i(\mathbf{r})}_{\mathbf{k}} - i k_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i} \\ &= -\hat{N}_{\mathbf{k},i} - i k_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i}, \quad \dots (82) \end{aligned}$$

$$k_i \hat{u}_{\mathbf{k},i} = 0.$$



Pressure computation

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u}, \quad \dots(80)$$

$$\nabla \cdot \mathbf{u} = 0, \quad \dots(81)$$

$$-\nabla^2 p = \nabla \cdot [\mathbf{u} \cdot \nabla \mathbf{u}] = \nabla \cdot \mathbf{N},$$

$$\hat{p}_{\mathbf{k}} = i \frac{1}{k^2} \mathbf{k} \cdot \hat{\mathbf{N}}_{\mathbf{k}},$$

$$\begin{aligned} \frac{d}{dt} \hat{u}_{\mathbf{k},i} &= -i k_j \widehat{u_j(\mathbf{r}) u_i(\mathbf{r})}_{\mathbf{k}} - i k_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i} \\ &= -\hat{N}_{\mathbf{k},i} - i k_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i}, \quad \dots(82) \end{aligned}$$

$$\hat{\mathbf{u}}_{\mathbf{k}}^{(n+1)} = \hat{\mathbf{u}}_{\mathbf{k}}^{(n)} [1 - \nu k^2 (\Delta t)] + (\Delta t) [-\hat{\mathbf{N}}_{\mathbf{k}} - i \mathbf{k} \hat{p}_{\mathbf{k}}^{(n)}]$$

$$\begin{aligned}\frac{d}{dt}\hat{u}_{\mathbf{k},i} &= -ik_j u_j(\mathbf{r})\widehat{u_i(\mathbf{r})}_{\mathbf{k}} - ik_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i} \\ &= -\hat{N}_{\mathbf{k},i} - ik_i \hat{p}_{\mathbf{k}} - \nu k^2 \hat{u}_{\mathbf{k},i}, \quad \dots(82)\end{aligned}$$

$$\Delta t < \tau_\nu = \frac{2h^2}{\nu\pi^2}.$$

$$\tau_{NL} = h / u_{rms}.$$

$$\tau_\nu / \tau_{NL} \approx u_{rms} h / \nu \approx Re / N,$$

$$\frac{\tau_\nu}{\tau_{NL}} \approx \frac{u_{rms} h}{\nu} \approx \frac{Re}{N} \approx \frac{Re}{Re^{3/4}} \approx Re^{1/4}.$$

Schrödinger Equation



$$i\hbar\partial_t\psi = \left[-\frac{\hbar^2}{2m}\nabla^2 + V(\mathbf{r}) \right] \psi ,$$

$$t = t'(\hbar/E_a) \text{ and } \mathbf{r} = \mathbf{r}'r_a,$$

$$i\partial_{t'}\psi = \left[-\frac{\alpha}{2}\nabla'^2 + V'(\mathbf{r}') \right] \psi ,$$

$$\alpha = \hbar^2/(mr_a^2E_a) \text{ and } V'=V/E_a.$$

equation and obtain

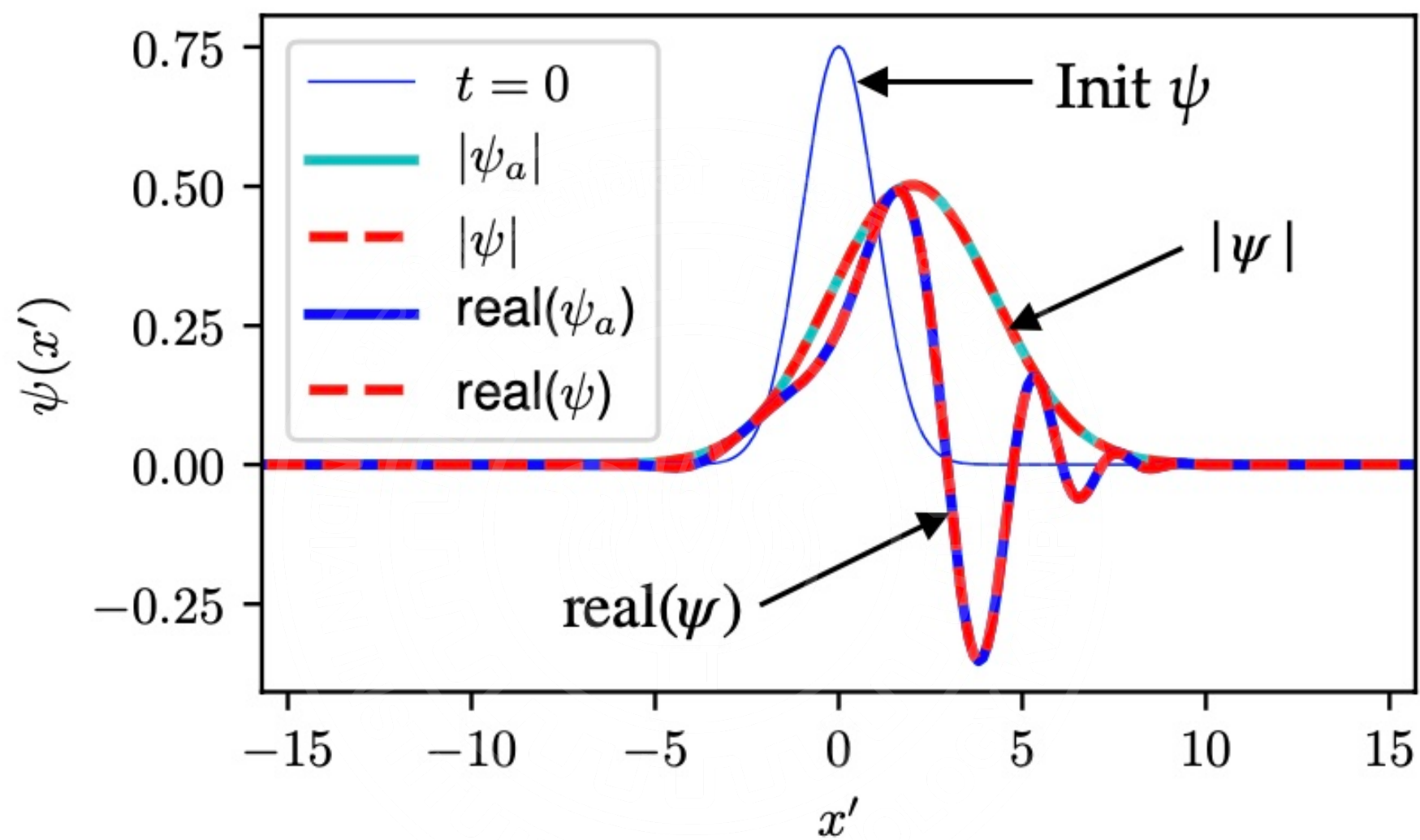
$$\partial_{t'}\hat{\psi}_{\mathbf{k}'} = -i\frac{\alpha}{2}k'^2\hat{\psi}_{\mathbf{k}'} - i\widehat{(V'\psi)}_{\mathbf{k}'} ,$$

$$\partial_{t'} \hat{\psi}_{\mathbf{k}'} = -i \frac{\alpha}{2} k'^2 \hat{\psi}_{\mathbf{k}'} - i \widehat{(V' \psi)}_{\mathbf{k}'},$$

$$\tilde{\psi}_{\mathbf{k}'} = \hat{\psi}_{\mathbf{k}'} \exp(i \frac{\alpha}{2} k'^2 t').$$

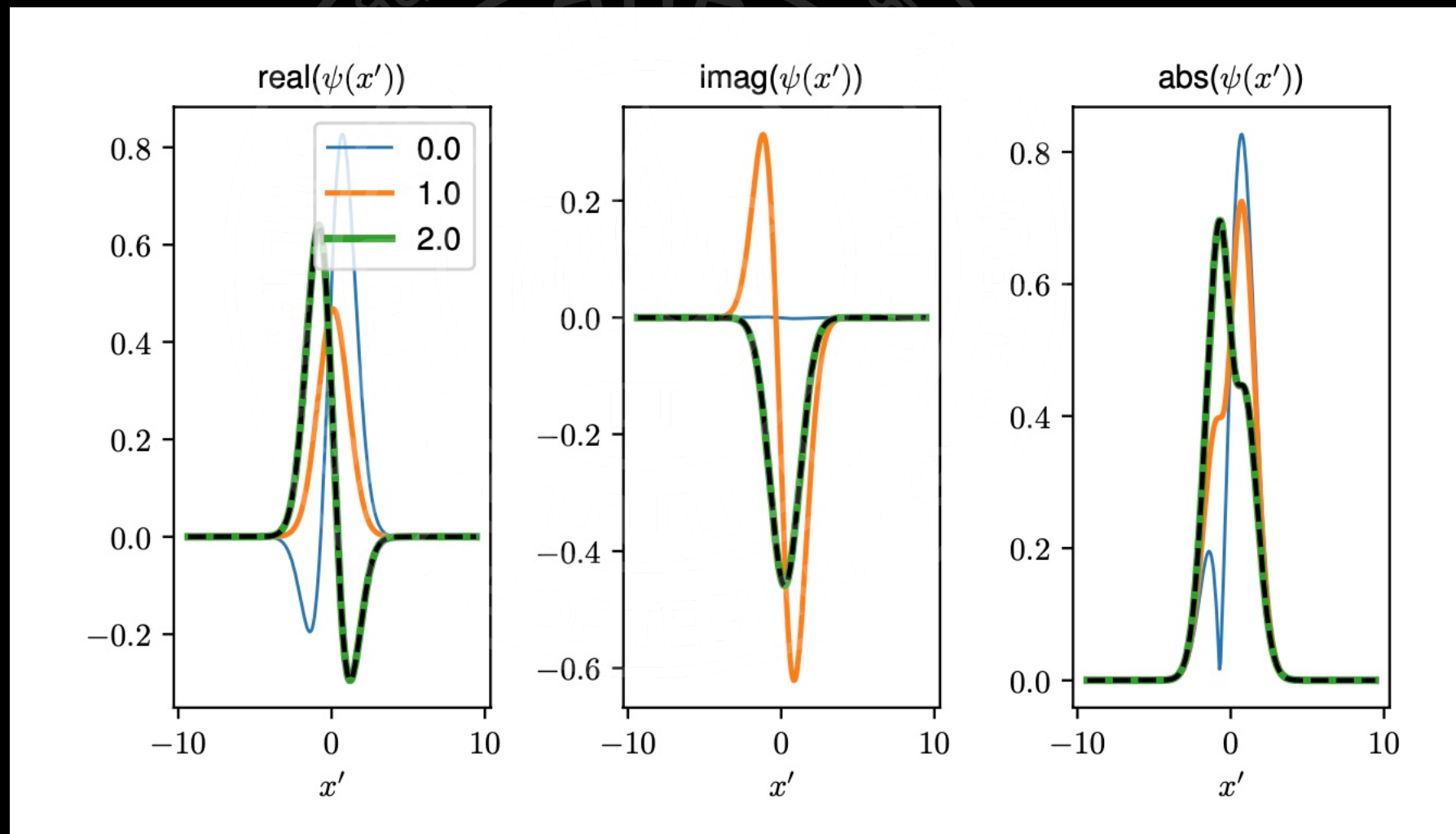
$$\partial_{t'} \tilde{\psi}_{\mathbf{k}'} = -i \widehat{(V' \psi)}_{\mathbf{k}'} \exp(i \frac{\alpha}{2} k'^2 t').$$

$$\hat{\psi}_{\mathbf{k}'}^{(n+1)} = \hat{\psi}_{\mathbf{k}'}^{(n-1)} \exp(-i \alpha k'^2 (\Delta t')) - i (2 \Delta t') \widehat{(V' \psi)}_{\mathbf{k}'} \exp(-i \frac{\alpha}{2} k'^2 (\Delta t'))$$



Quantum Oscillator

$$i\partial_t\psi = -\frac{1}{2}\frac{d^2\psi}{dx'^2} + \frac{1}{2}x'^2\psi.$$



Gross–Pitaevskii Equation

$$i\hbar\partial_t\psi = \left[-\frac{\hbar^2}{2m}\nabla^2 + V(\mathbf{r}) + g|\psi|^2 \right] \psi ,$$

$$i\hbar\partial_t\hat{\psi}_{\mathbf{k}} = \frac{\hbar^2 k^2}{2m}\hat{\psi}_{\mathbf{k}} + \widehat{(V\psi)}_{\mathbf{k}} + (g\widehat{|\psi|^2}\psi)_{\mathbf{k}} ,$$

